

An exceed 60% efficiency Nd:YAG transparent ceramic laser with low attenuation loss effect

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Keywords: attenuation loss, side-pumped, end-pumped, Nd:YAG ceramic laser, optical conversion efficiency.

15 Abstract

16 Here, attenuation loss effect and laser performance enhancement of Nd:YAG transparent ceramics were investigated. Using a 0.6 at.% Nd:YAG ceramic rod with 3 mm diameter and 65 mm length, the 17 scattering coefficient and absorption coefficient at 1064 nm were measured to be 0.0001 cm⁻¹ and 18 19 0.0017 cm⁻¹, respectively. For the 808 nm side-pumped laser experiment, an average output power of 44.9 W was achieved with an optical-to-optical conversion efficiency of 26.4%, which was nearly 20 21 same with that of 1 at.% single crystal. Adopting the 885 nm direct end-pumped scheme, the 22 following laser tests demonstrated the high optical efficiency of 62.5% and maximum output power of 144.8 W were obtained at absorbed pump power of 231.5 W. This was up to now the highest 23 optical conversion efficiency acquired in Nd:YAG ceramic laser to our knowledge. It proves that 24 25 high power and high efficiency laser output could be generated by high optical quality Nd:YAG ceramic rod along with the 885 nm direct pumping technology. 26

27 1 Introduction

28 Polycrystalline transparent ceramic materials have become an attractive alternative to widely used 29 single-crystals because of their favorable characteristics, such as higher doping concentration, larger 30 scale, more function design freedom, easier manufacture, low cost, and especially superior resistance to fracture [1]. Since an effective laser output with polycrystalline Nd:YAG ceramics was first 31 32 performed in 1995 [2], numerous attempts have been delivered on the field of high power and high 33 efficiency Nd:YAG ceramic solid-state lasers. For example, these include the output power breaking 34 the 1 kW mark in 2002 and then the remarkable demonstration of more than 100 kW from a YAG 35 ceramic laser system in 2009 [3]. And for middle and high power laser oscillation, the increased optical conversion efficiencies from 14.5% to 52.5% have been reported one by one [4-7]. Among the significantly milestone achievements, the gain medium with high optical quality is the key factor for highly efficient laser oscillation. Therefore, optical properties including optical absorption, emission spectra and fluorescence lifetime have been widely studied for Nd:YAG ceramics [8-10], and very similar results were obtained with that of Nd:YAG single crystals. However, the wellknown attenuation loss that has enormous influence on the laser performance is rarely available.

42 Attenuation loss mainly include the scattering and absorption effect caused by the residual pores, grain boundary phases, secondary phases and impurity ions, which will prohibit laser output. In 1998, 43 44 Ikesue et al. showed the scattering coefficients of Nd:YAG ceramics obtained from Fresnel's 45 equation by an optical spectra method [11]. Subsequently, Li et al. presented the absorption 46 coefficients of Nd:YAG ceramics at the laser wavelength [12]. Unfortunately, they did not gave a 47 precise distinction between the scattering coefficients and the absorption coefficients. Recently, 48 Boulesteix et al. computed the light scattering of Nd:YAG ceramics according to Mie light scattering 49 theory. However, the absorption effect was ignored [13]. Providing accurate measurement of the scattering and absorption coefficient could promote the materials research, such as laser experiment 50 51 and fabrication technology.

52 This paper introduced an effective measured method of light attenuation loss by means of an 53 integrating sphere technique, which is used to analyze the optical properties of Nd:YAG ceramic and crystal. A 0.6 at.%-doped Nd:YAG ceramic sample has well optical quality with the scattering 54 coefficient of 0.0001 cm⁻¹ and absorption coefficient of 0.0017 cm⁻¹. Here, an 808 nm laser diode 55 (LD) side-pumped laser configuration was developed, delivering a 26.4% of optical efficiency with 56 57 44.9 W of output power. Moreover, an end-pumped linear cavity with an 885 nm LD pumping was 58 designed to improve the optical conversion efficiency. As a result, a maximum output power of 144.8 59 W was obtained under the absorbed pump power of 231.5 W. The corresponding optical efficiency is calculated to be as high as 62.5%, which is a significant improvement for high efficiency Nd:YAG 60 ceramic lasers. 61

62 2 Absorption and scattering coefficient measurement



63 64

Figure 1.Experiment schematic for the scattering coefficient measurements

To assess the overall optical quality of the material, the measurements of light scattering and absorption were carried out at 1064 nm, based on a homemade scattering loss analyzer with an integrating sphere. The measurement configuration was as shown in Figure 1. When a laser beam is nearly at normal incidence upon the samples mounted in the center of integrating sphere, part of the radiation is reflected, part is scattered, part is absorbed, and the rest is transmitted. The incident laser power was denoted to be P_{in} , and the scattering power and absorption power in the material were

- 71 respectively described as P_s and P_a . The transmission power passing through the samples was defined
- 72 as P_T , and the Fresnel reflection coefficient of the front and back surface of the samples was denoted
- to be *r*. According to the law of Fresnel reflection and Lambert-Beer [14], the correlated relationship
- 74 of above power distributions can be written as

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$$\begin{cases} P_T = (P_{in} - rP_{in} - P_a - P_s)(1 - r) \\ P_a + P_s = P_{in}(1 - r)[1 - \exp(-\alpha L)] \end{cases}$$
 (1)

where α is the attenuation coefficient of the sample, and *L* is the length of the sample. It is evident that the power attenuation loss inside the sample is caused by the scattering and absorption mechanism. Thus, we could obtain $\alpha = \alpha_a + \alpha_s$, and $P_{\alpha'}/P_s = \alpha_{\alpha'}/\alpha_s$, where α_a and α_s are the absorption coefficient and the scattering coefficient, respectively. After simplifying equation (1), yields

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$$\begin{cases}
\frac{P_T}{P_{in}} = (1-r)^2 - \frac{P_s}{P_{in}} (\frac{\alpha_a}{\alpha_s} + 1)(1-r) \\
\frac{P_s}{P_{in}} (\frac{\alpha_a}{\alpha_s} + 1) = (1-r) \{1 - \exp[-(\alpha_a + \alpha_s)L]\}
\end{cases}$$
(2)

81 While P_{in} and P_T were measured by a power meter PM (NOVA II OPHIR). The scattered light 82 intensity with and without the samples was collected with a calibrated photoelectric detector PD 83 (Thorlabs Inc., DET200) mounted at the top of the integrating sphere, respectively recording as P_1 84 and P_2 . Here, the value of P_s/P_{in} is equal to the ratio of P_1/P_2 . By solving equation (2), the absorption 85 coefficient α_a and the scattering coefficient α_s could be achieved.

In the measurement, two Nd:YAG ceramic samples with 0.6 at.% and 1.0 at.% doping concentration, and a Nd:YAG single crystal with 1.0 at.% doping concentration were employed, which were fabricated by Nanyang Technological University. Each sample has a size of 3 mm diameter and 65 mm length, and both facets of samples are polished and antireflection coated at 1064 nm to reduce the surface reflection. Therefore, the reflectivity r at the surface of the sample is assumed to be about 0.1%.

Parameters	Samples		
	S1	S2	S 3
Nd ³⁺ doped concentration (at. %)	0.6	1.0	1.0
Ratio of transmittance $(P_T/P_{in}, \%)$	0.986	0.981	0.987
Ratio of scattering $(P_{s}/P_{in}, \%)$	0.0007	0.009	0.007
Scattering coefficient (α_s , cm ⁻¹)	0.0001	0.0014	0.0011
Absorption coefficient (α_a , cm ⁻¹)	0.0017	0.0012	0.0006
Attenuation coefficient (α , cm ⁻¹)	0.0018	0.0026	0.0017

92 **Table 1.** Measured values and correlative results at 1064 nm of different samples

Table 1 summarizes the corresponding scattering coefficient and absorption coefficient of each sample. Obviously, the crystal sample has the highest optical quality with the smallest attenuation coefficient of 0.0017 cm⁻¹, where the scattering coefficient and absorption coefficient were measured to be 0.0011 cm⁻¹ and 0.0006 cm⁻¹. Compared with 1 at.% Nd:YAG ceramic rod, the 0.6 at.% Nd:YAG ceramic rod with attenuation loss of 0.0018 cm⁻¹ is nearly the same as single crystal, which

98 could be easier to produce high power laser output. In addition, the existence of low impurity ions 99 during preparation process is inevitable, which results in a large absorption coefficient with a same 100 order of magnitude as the scattering coefficient and could not be neglected in the defects of the 101 ceramic materials. The above data indicate that ceramic YAG is essentially identical for single crystal 102 YAG in optical properties measured, especially for the low scattering and absorption losses.

103 **3 Laser Experiment**

104 In order to evaluate the laser performance of the ceramic samples compared with the Nd:YAG 105 crystal, a compact flat-flat cavity was adopted and side-pumped by LD at the wavelength of 808 nm 106 for high pump absorption efficiency. The samples were surrounded by arrays of diode lasers with the 107 total pump power of 180 W. The Nd:YAG rod and LD arrays were cooled to be 25°C by deionized water flowing, to match the pump radiation wavelength of LD and the 808.5 nm absorption spectrum 108 109 of Nd:YAG. The mirror M1 was coated with high reflectance (HR) at 1064 nm, and the mirror M2 110 was an output coupler with partially-reflectivity of 80% at 1064 nm. The 1064 nm output power was 111 monitored by the PM. The total cavity length is about 70 mm.



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Figure 2.Output power of 1064 nm laser versus pump power at 808 nm.

114 The average laser output power of three Nd:YAG samples as a function of LD pump power at 115 808 nm is shown in Figure 2. The output power grows approximately linearly with the increase of 116 pump power, and does not show any roll-over effect, indicating that higher output can be achieved 117 with increasing pump energy continuously. With 170 W maximum pump power, 44.9 W and 46.2 W 118 laser output were obtained at 1064 nm for 0.6 at.%Nd:YAG ceramic and1 at.% Nd:YAG crystal, 119 respectively. The corresponding optical-to-optical conversion efficiencies are 26.4% and 27.2%. The 120 optical efficiency of the 0.6 at.% ceramic laser is only 0.8 % less than that of the single crystal laser, 121 due to the difference of absorbed pump power caused by different neodymium concentration. It was 122 proven, from the aspect of output power and laser efficiency, the ceramic and crystal materials share 123 almost same laser characteristics. For 1 at.%-doped ceramic sample, laser output of 38.6 W was 124 lower than the same doping concentration crystal because of the largest attenuation loss coefficient 125 and serious thermal effect.

The quantum defect between the pump and laser emission wavelength is one of the major factors that limit the LD pumped solid-state lasers to generate high power and high efficiency [15]. Compared to traditional 808 nm pumping, adopting 885 nm diodes will have a reduction of thermal load nearly by 30% and will thus lead to an improvement in the overall laser efficiency. The 0.6 at.%Nd:YAG ceramic rod was employed as the measured sample, and an end-pumped plane-plane
linear cavity was designed. An 885 nm fiber-coupled diode laser (DILAS, 400 µm diameter and 0.22
NA) was used as the pump source, delivering themaximum power of 250 W. It is focused into the
ceramic sample by the coupling lens of 1:1. The laser sample was cooled by the re-circulating filtered
water at 16°C. The input mirror M1 was coated with high transmission film at pump wavelength of
885 nm and HR film at 1064nm. And, the output coupler M2 has a transmission of 20% at 1064 nm.
The cavity length of the resonator is about 73 mm to keep the cavity mode in the sample matching

the pump mode. The mirror M3 is adopted to separate pump light and output laser.



Figure 3. (a) Absorbed power versus incident pump power at 885 nm. (b) Output power of 1064 nm laser versus absorbed power at 885 nm.

141 The optical efficiency for a reasonable comparison could be calculated based on absorbed pump 142 power. Firstly, absorbed pump power was estimated by monitoring the pump power passing through the sample at different incident levels, as displayed in Figure 3(a). The absorption power increased 143 and the absorption coefficient decreased with increasing the input pump power. For instance, the 144 absorption coefficient varied from 0.6 cm⁻¹ to 0.43 cm⁻¹, corresponds to a pump absorption of 97.9% 145 146 and 94%, which is attributed to the absorption saturation behavior of lower doping concentration. As 147 shown in Figure 3(b), the output power at 1064 nm increased linearly in accordance with the 148 absorbed pump power. At the absorbed pump power of 231.5 W, the maximum output power was as 149 high as 144.8 W with a corresponding optical-to-optical conversion efficiency of 62.5%. The optical 150 conversion efficiency versus absorbed pump power is also given in Figure 3(b). Actually, the 151 maximum conversion efficiency of about 64.6% was obtained at 205 W absorbed pumping. To the 152 best of our knowledge, this is the highest optical conversion efficiency of all 1064 nm laser systems 153 with end pumped laser modules. The drop in efficiency after 205 W was caused by the serious 154 thermal effect in laser ceramic rod for high power operation.

155 4 Discussion and Conclusion

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In conclusion, the comparison of laser performance of Nd:YAG ceramics and crystal as well as attenuation loss are introduced and analyzed, based on an integrating sphere and 808 nm LD sidepumped laser experiment. As a result, the Nd:YAG ceramic could be processed to access almost identical optical properties with the single crystal. Moreover, a 0.6 at.%Nd:YAG ceramic rod was further investigated for producing high optical conversion efficiency, by means of 885 nm LD direct end-pumped technology. Under the absorbed power of 231.5 W, the maximum output power of 144.8

- 162 W was obtained with an optical efficiency of 62.5%, which is the highest efficiency 1064 nm
- 163 Nd:YAG ceramic laser ever reported.

164 **5** Conflict of Interest

165 The authors declare that the research was conducted in the absence of any commercial or financial166 relationships that could be construed as a potential conflict of interest.

167 **6** Author Contributions

168 Yu Shen conceived the project. Jin-Quan Chang conducted the experiment. Qi Bian wrote the

- 169 manuscript and all authors contributed to discussions during its preparation. Yong Bo and Qin-Jun
- 170 Peng supervised the project.

171 **7 Funding**

172 This work was supported by Supported by the National Key Research and Development Program of

173 China(key special projects, Young Scientist Program, No. 2022YFB3607900), Key Laboratory 174 Foundation of Chinese Academy of Sciences, Key Lab of Solid State Laser(No. CXJJ-22S020), the

175 National Science Foundation for Young Scientists of China (Grant No. 11504389, 51890864,

175 National Science Foundation for Found Sciencists of China (Grant No. 11504589, 51890804;

62005295). The authors acknowledge some crystal support from Chengdu Dien PhotoelectricTechnology Co., Ltd. (DIEN TECH).

178 **8** Acknowledge

179 The authors acknowledge some crystal support from Chengdu Dien Photoelectric Technology Co.,180 Ltd. (DIEN TECH).

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